

Quantum Field Theory

Problem Sheet 7

Problem 15: Quantization of Maxwell Field

The operators of the transversal vector potential $\hat{A}_{\perp i}(\mathbf{x}, t)$ and their canonical conjugated momenta $\hat{\pi}_i(\mathbf{x}, t) = \epsilon_0 \frac{\partial \hat{A}_{\perp i}(\mathbf{x}, t)}{\partial t}$ in the Heisenberg picture obey the equal-time commutation relations

$$\left[\hat{\pi}_i(\mathbf{x}, t), \hat{\pi}_j(\mathbf{x}', t) \right]_- = \left[\hat{A}_{\perp i}(\mathbf{x}, t), \hat{A}_{\perp j}(\mathbf{x}', t) \right]_- = 0, \quad \left[\hat{A}_{\perp i}(\mathbf{x}, t), \hat{\pi}_j(\mathbf{x}', t) \right]_- = i\hbar \delta_{ij}^{\perp}(\mathbf{x} - \mathbf{x}') \quad (1)$$

with the transversal delta function

$$\delta_{ij}^{\perp}(\mathbf{x} - \mathbf{x}') = \delta_{ij} \delta(\mathbf{x} - \mathbf{x}') + \frac{1}{4\pi} \partial'_j \partial'_i \frac{1}{|\mathbf{x} - \mathbf{x}'|}. \quad (2)$$

a) Demonstrate that the commutation relations (1) are compatible with the quantized version of the transversality condition

$$\partial_i \hat{A}_{\perp i}(\mathbf{x}, t) = 0. \quad (3)$$

Here we have used the Einstein summation convention that implies a summation over equal indices. (2 points)

b) In the lecture it is shown that the field operator $\hat{\mathbf{A}}_{\perp}(\mathbf{x}, t)$ has the following Fourier decomposition:

$$\hat{\mathbf{A}}_{\perp}(\mathbf{x}, t) = \sum_{\lambda=\pm} \int d^3k N_{\mathbf{k}} \left[\boldsymbol{\epsilon}(\mathbf{k}, \lambda) e^{i(\mathbf{k}\mathbf{x} - \omega_{\mathbf{k}}t)} \hat{a}_{\mathbf{k}, \lambda} + \boldsymbol{\epsilon}(\mathbf{k}, \lambda)^* e^{-i(\mathbf{k}\mathbf{x} - \omega_{\mathbf{k}}t)} \hat{a}_{\mathbf{k}, \lambda}^{\dagger} \right] \quad (4)$$

Here the dispersion reads $\omega_{\mathbf{k}} = c|\mathbf{k}|$ and the polarization vectors $\boldsymbol{\epsilon}(\mathbf{k}, \lambda)$ fulfill the properties

$$\mathbf{k} \boldsymbol{\epsilon}(\mathbf{k}, \lambda) = 0, \quad \boldsymbol{\epsilon}(\mathbf{k}, \lambda) \boldsymbol{\epsilon}(\mathbf{k}, \lambda')^* = \delta_{\lambda, \lambda'}, \quad \boldsymbol{\epsilon}(-\mathbf{k}, \lambda) = \boldsymbol{\epsilon}(\mathbf{k}, -\lambda) = \boldsymbol{\epsilon}(\mathbf{k}, \lambda)^*. \quad (5)$$

Prove

$$\hat{a}_{\mathbf{k}, \lambda} = \frac{1}{2(2\pi)^3 N_{\mathbf{k}}} \int d^3x \boldsymbol{\epsilon}(-\mathbf{k}, \lambda) e^{-i(\mathbf{k}\mathbf{x} - \omega_{\mathbf{k}}t)} \left[\hat{\mathbf{A}}_{\perp}(\mathbf{x}, t) + \frac{i}{\epsilon_0 \omega_{\mathbf{k}}} \hat{\boldsymbol{\pi}}(\mathbf{x}, t) \right] \quad (6)$$

and the corresponding relation for $\hat{a}_{\mathbf{k}, \lambda}^{\dagger}$. (4 points)

c) Determine the commutators

$$\left[\hat{a}_{\mathbf{k}, \lambda}, \hat{a}_{\mathbf{k}', \lambda'} \right]_- = ?, \quad \left[\hat{a}_{\mathbf{k}, \lambda}^{\dagger}, \hat{a}_{\mathbf{k}', \lambda'}^{\dagger} \right]_- = ?, \quad \left[\hat{a}_{\mathbf{k}, \lambda}, \hat{a}_{\mathbf{k}', \lambda'}^{\dagger} \right]_- = ?. \quad (7)$$

How do you have to fix the normalization constants $N_{\mathbf{k}}$, so that the operators $\hat{a}_{\mathbf{k},\lambda}$ and $\hat{a}_{\mathbf{k},\lambda}^\dagger$ can be interpreted as annihilation and creation operators of a single particle? (6 points)

d) According to the lecture the Hamilton operator of the electromagnetic field reads

$$\hat{H} = \frac{1}{2} \int d^3x \left[\frac{1}{\epsilon_0} \hat{\pi}_i(\mathbf{x}, t) \hat{\pi}_i(\mathbf{x}, t) + \frac{1}{\mu_0} \partial_i \hat{A}_{\perp j}(\mathbf{x}, t) \partial_i \hat{A}_{\perp j}(\mathbf{x}, t) \right]. \quad (8)$$

Show that the normal ordered Hamilton operator $: \hat{H} := \hat{H} - \langle 0 | \hat{H} | 0 \rangle$ has the following representation:

$$: \hat{H} := \sum_{\lambda=\pm 1} \int d^3k \hbar \omega_{\mathbf{k}} \hat{a}_{\mathbf{k},\lambda}^\dagger \hat{a}_{\mathbf{k},\lambda}. \quad (9)$$

(4 points)

e) The momentum operator of the electromagnetic field reads

$$\hat{\mathbf{p}} = \int d^3x \left[\nabla \times \hat{\mathbf{A}}_{\perp}(\mathbf{x}, t) \right] \times \hat{\boldsymbol{\pi}}(\mathbf{x}, t). \quad (10)$$

Demonstrate that the normal ordered momentum operator $: \hat{\mathbf{p}} := \hat{\mathbf{p}} - \langle 0 | \hat{\mathbf{p}} | 0 \rangle$ has the following representation

$$: \hat{\mathbf{p}} := \sum_{\lambda=\pm 1} \int d^3k \hbar \mathbf{k} \hat{a}_{\mathbf{k},\lambda}^\dagger \hat{a}_{\mathbf{k},\lambda}. \quad (11)$$

(4 points)

f) The spin operator of the electromagnetic field reads

$$\hat{\mathbf{S}} = \int d^3x \hat{\mathbf{A}}_{\perp}(\mathbf{x}, t) \times \hat{\boldsymbol{\pi}}(\mathbf{x}, t). \quad (12)$$

Prove that the normal ordered spin operator $: \hat{\mathbf{S}} := \hat{\mathbf{S}} - \langle 0 | \hat{\mathbf{S}} | 0 \rangle$ possesses the following representation

$$: \hat{\mathbf{S}} := \sum_{\lambda=\pm 1} \int d^3k \lambda \hbar \frac{\mathbf{k}}{|\mathbf{k}|} \hat{a}_{\mathbf{k},\lambda}^\dagger \hat{a}_{\mathbf{k},\lambda}. \quad (13)$$

Hint: Use the following identity for the polarization vectors

$$\boldsymbol{\epsilon}(\mathbf{k}, \lambda) \times \boldsymbol{\epsilon}(\mathbf{k}, \lambda')^* = -i \lambda \frac{\mathbf{k}}{|\mathbf{k}|} \delta_{\lambda \lambda'}. \quad (14)$$

(4 points)

Drop the solutions in the post box on the 5th floor of building 46 or, in case of illness/quarantine, send them via email to jkrauss@rhrk.uni-kl.de until June 14 at 14.00.