

neutrinos without rest mass:

$$\psi = \begin{pmatrix} \chi \\ \xi \end{pmatrix} \rightarrow \frac{1}{2} (1 - \gamma^5) \psi = \begin{pmatrix} \chi \\ 0 \end{pmatrix}$$

$$\rightarrow \frac{1}{2} (1 + \gamma^5) \psi = \begin{pmatrix} 0 \\ \xi \end{pmatrix}$$

chirality operator $\gamma^5 = \begin{pmatrix} -\mathbb{I} & 0 \\ 0 & \mathbb{I} \end{pmatrix}$

$$\gamma^5 \frac{1}{2} (1 \mp \gamma^5) \psi = \mp \frac{1}{2} (1 \mp \gamma^5) \psi$$

chirality

action: $\Delta = \int \bar{\psi} i \gamma^5 \partial_\mu \frac{1}{2} (1 \mp \gamma^5) \psi$

no rest mass term
mixing χ and ξ

Weyl equation:

$$\frac{\delta \Delta}{\delta \bar{\psi}(x)} = \frac{\partial \mathcal{L}}{\partial \bar{\psi}(x)} - \partial_\mu \frac{\partial \mathcal{L}}{\partial \partial_\mu \bar{\psi}(x)} = 0 \Rightarrow i \gamma^\mu \partial_\mu \frac{1}{2} (1 \mp \gamma^5) \psi = 0$$

not to solve

fixed momentum: $\psi(x) = \psi e^{-i p \cdot x / \hbar}$

$$\Rightarrow \vec{\sigma} \cdot \vec{p} \frac{1}{2} (1 \mp \gamma^5) \psi = \gamma^0 p^0 \frac{1}{2} (1 \mp \gamma^5) \psi \quad | \quad \gamma^5 \gamma^0$$

$$1) \gamma^5 \gamma^0 \gamma^k = \begin{pmatrix} -\mathbb{I} & 0 \\ 0 & \mathbb{I} \end{pmatrix} \begin{pmatrix} 0 & \mathbb{I} \\ \mathbb{I} & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^k \\ -\sigma^k & 0 \end{pmatrix} = \begin{pmatrix} \sigma^k & 0 \\ 0 & \sigma^k \end{pmatrix}$$

$$2) \gamma^5 (\gamma^0 \gamma^0) = \gamma^5$$

$$\Rightarrow \underbrace{\frac{1}{2} \begin{pmatrix} \sigma^k & 0 \\ 0 & \sigma^k \end{pmatrix} p^k}_{= S^k} \frac{1}{2} (1 \mp \gamma^5) \psi = \frac{1}{2} p^0 \underbrace{\gamma^5}_{\text{sgn}(p^0)} \frac{1}{2} (1 \mp \gamma^5) \psi$$

$\vec{S} \cdot \vec{p}$

$$\underbrace{\frac{\vec{S} \cdot \vec{P}}{|\vec{P}|}}_{\text{helicity operator}} \frac{1}{2} (1 \mp \gamma^5) \psi = \frac{1}{2} (\text{sgn } p_0) \underbrace{\gamma^5 \frac{1}{2} (1 \mp \gamma^5) \psi}_{\mp \frac{1}{2} (1 \mp \gamma^5) \psi} = \mp \frac{1}{2} \text{sgn } p_0$$

helicity = chirality for these massless particles

Note: For massive spin 1/2 particle helicity and chirality are different quantities

6.10 Charge Conjugation:

investigate another discrete symmetry

$$\psi(x) \longrightarrow \psi^*(x)$$

$$\psi_c^{\dagger}(x) = C \underbrace{\bar{\psi}^T(x)}_{=?} = C \gamma^0 \psi^*(x)$$

$$\bar{\psi}(x) = \psi^{\dagger}(x) \gamma^0, \quad \bar{\psi}^T(x) = (\gamma^0)^T (\psi^{\dagger}(x))^T = \gamma^0 \psi^*(x)$$

$$\left(i \gamma^{\mu} \partial_{\mu} - m \right) \psi(x) = 0 \quad \longleftrightarrow \quad \left(i \gamma^{\mu} \partial_{\mu} - m \right) \psi_c^{\dagger}(x) = 0$$

\downarrow \downarrow

$$\begin{aligned}
 & -i \partial_\mu \psi(x) \gamma^\mu - m \psi(x) = 0 \quad \left. \begin{array}{l} \\ \\ \end{array} \right\} \cdot \delta^0 \quad (i \gamma^\mu \partial_\mu - m) C \bar{\psi}^T(x) = 0 \quad | \cdot C^{-1} \\
 & \downarrow \psi(x) = \bar{\psi}(x) \delta^0 \\
 & -i \partial_\mu \bar{\psi}(x) \gamma^0 \gamma^\mu - m \bar{\psi}(x) \delta^0 = 0 \quad \left. \begin{array}{l} \\ \\ \end{array} \right\} \cdot \delta^0 \quad (i C^{-1} \gamma^\mu C \partial_\mu - C^{-1} m) \bar{\psi}^T(x) = 0 \quad | \cdot T \\
 & i \partial_\mu \bar{\psi}(x) \delta^0 \gamma^{\mu\dagger} \delta^0 + m \bar{\psi}(x) = 0 \quad \underline{i \partial_\mu \bar{\psi}(x) (C^{-1} \gamma^\mu C)^T - m \bar{\psi}(x) = 0}
 \end{aligned}$$

$$= \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & \gamma^0 \\ \gamma^0 & 0 \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} = \begin{pmatrix} 0 & \gamma^0 \\ \gamma^0 & 0 \end{pmatrix} = \gamma^0$$

$$\underline{i \partial_\mu \bar{\psi}(x) \gamma^\mu + m \bar{\psi}(x) = 0}$$

$$\begin{aligned}
 & (C^{-1} \gamma^\mu C)^T = -\gamma^\mu \\
 & C^{-1} \gamma^\mu C = -(\gamma^\mu)^T \\
 & \text{to be solved for } C
 \end{aligned}$$

Solution ansatz: $C = \begin{pmatrix} c & 0 \\ 0 & -c \end{pmatrix}$, $C^{-1} = \begin{pmatrix} c^{-1} & 0 \\ 0 & -c^{-1} \end{pmatrix}$
 c : 2×2 -matrix

$$\begin{aligned}
 & \begin{pmatrix} c^{-1} & 0 \\ 0 & -c^{-1} \end{pmatrix} \begin{pmatrix} 0 & \gamma^0 \\ \gamma^0 & 0 \end{pmatrix} \begin{pmatrix} c & 0 \\ 0 & -c \end{pmatrix} = \dots = \begin{pmatrix} 0 & -c^{-1} \gamma^0 c \\ -c^{-1} \gamma^0 c & 0 \end{pmatrix} \\
 & = - \begin{pmatrix} 0 & \gamma^0 \\ \gamma^0 & 0 \end{pmatrix}^T = - \begin{pmatrix} 0 & (\gamma^0)^T \\ (\gamma^0)^T & 0 \end{pmatrix} \rightarrow \begin{cases} c^{-1} \gamma^0 c = +(\gamma^0)^T \\ c^{-1} \gamma^0 c = +(\gamma^0)^T \end{cases}
 \end{aligned}$$

$\mu=0$: $c^{-1} \gamma^0 c = (\gamma^0)^T$, $\mu=1$: $c^{-1} \gamma^1 c = -(\gamma^1)^T$

Pauli matrices:

$$\gamma^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, (\gamma^0)^T = \gamma^0; \quad \gamma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, (\gamma^1)^T = \gamma^1$$

$$\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, (\sigma^2)^T = -\sigma^2; \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, (\sigma^3)^T = \sigma^3$$

Ansatz: $C = -i\sigma^2 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$

$$C^+ = C^{-1} = C^T = -C = -C^*$$

$$C^{-1}\sigma^0 C = i\sigma^2\sigma^0(-i\sigma^2) = \sigma^2\sigma^0\sigma^2 = (\sigma^2)^2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \sigma^0 = (\sigma^0)^T \quad \mu=0$$

$$C^{-1}\sigma^1 C = i\sigma^2\sigma^1(-i\sigma^2) = \sigma^2\sigma^1\sigma^2 = -\sigma^1(\sigma^2)^2 = -\sigma^1 = -(\sigma^1)^T \quad \mu=1$$

\Rightarrow show also $\mu=2, \mu=3$

Conclusion: $C^+ = C^{-1} = C^{-2} = -C = -C^*$

representation of C in δ -matrices:

$$i\gamma^0\gamma^2 = i \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^2 \\ -\sigma^2 & 0 \end{pmatrix} = \begin{pmatrix} C & 0 \\ 0 & -C \end{pmatrix} = C$$

Involutive properties:

$$\psi_c^\dagger(x) = C\gamma^0\psi_c^*(x) = C\gamma^0[C\gamma^0\psi^*(x)]^* = C\gamma^0 C^* \gamma^0 \psi(x)$$

$$= \begin{pmatrix} C & 0 \\ 0 & -C \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} C^* & 0 \\ 0 & -C^* \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} = \begin{pmatrix} -CC^* & 0 \\ 0 & -CC^* \end{pmatrix} \psi(x)$$

$$= \psi(x) \quad \checkmark$$

$$= \begin{pmatrix} 0 & C \\ -C & 0 \end{pmatrix} = \begin{pmatrix} 0 & C^* \\ -C^* & 0 \end{pmatrix}$$

$$C^* = -C^{-1}$$

Continuity equation:

$$i \partial^\mu \partial_\mu \psi(x) - m \psi(x) = 0 \xrightarrow{\bar{\psi}(x)} i \bar{\psi}(x) \partial^\mu \partial_\mu \psi(x) - m \bar{\psi}(x) \psi(x) = 0$$

$$i \partial_\mu \bar{\psi}(x) \partial^\mu \psi(x) + m \bar{\psi}(x) \psi(x) = 0 \xrightarrow{|\cdot \psi(x)} i \partial_\mu \bar{\psi}(x) \partial^\mu \psi(x) + m \bar{\psi}(x) \psi(x) = 0$$

$$\Rightarrow \bar{\psi}(x) \partial^\mu \partial_\mu \psi(x) + \partial_\mu \bar{\psi}(x) \partial^\mu \psi(x) = \partial_\mu \left[\bar{\psi}(x) \partial^\mu \psi(x) \right] = 0$$

current density: $\tilde{j}^\mu(x) = k \bar{\psi}(x) \gamma^\mu \psi(x) \sim \tilde{j}^\mu(x)$

conserved charge: $Q = \int d^3x \tilde{j}^0(\vec{x}, t) = k \int d^3x \psi^\dagger(x) \gamma^0 \gamma^0 \psi(x)$
 $= \psi^\dagger(x) \underbrace{\gamma^0 \gamma^0}_{=1} \psi(x)$

$$\Rightarrow Q = k \int d^3x \psi^\dagger(\vec{x}, t) \psi(\vec{x}, t)$$

$$\bar{\psi}'_C(x) = \psi^\dagger_C(x) \gamma^0 = (C \gamma^0 \psi^*(x))^\dagger \gamma^0 = \psi^\dagger(x) \underbrace{(\gamma^0)^\dagger}_{=\gamma^0} + \underbrace{C^\dagger}_{=-C} \gamma^0$$

$$= \psi^\dagger(x) C$$

$$= -\gamma^0 C \gamma^0 = \dots = C$$

charge conjugation of current density:

$$\tilde{j}'_C^\mu(x) = k \bar{\psi}'_C(x) \gamma^\mu \psi'_C(x) = k \psi^\dagger(x) C \gamma^\mu C \gamma^0 \psi^*(x)$$

$$= \left(\tilde{j}'_C^\mu(x) \right)^T = k \underbrace{\psi^\dagger(x)}_{=\bar{\psi}(x)} \underbrace{(\gamma^0)^\dagger}_{=\gamma^0} \underbrace{(C \gamma^\mu C)^T}_{\gamma^\mu} \psi(x) = \tilde{j}^\mu(x)$$

↑ C number no sign change

Note: We do not get the expected sign change

⇒ Revisit charge conjugation on level of 2nd quantization yields sign change.

6.11 Time Inversion: "run a movie backwards in time"

$$x = (x^0, \vec{x}) \rightarrow x'_T = T x = -\tilde{x} = (-x^0, \vec{x})$$

$$T = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad T^2 = 1 \quad (\Rightarrow) \quad T^{-1} = T$$

$$\text{show: } T^{-1} L_{\mu} T^{-1} = L_{\mu} \quad (\Rightarrow) \quad [T, L_{\mu}]_- = 0$$

$$T^{-1} M_{\mu\nu} T^{-1} = -M_{\mu\nu} \quad (\Rightarrow) \quad [T, M_{\mu\nu}]_+ = 0$$

Side consideration: Schrödinger

$$\left(i\hbar \frac{\partial}{\partial t} + \frac{\hbar^2}{2m} \Delta \right) \psi(\vec{x}, t) = 0$$

time-inverted wave function $\psi'_T(\vec{x}, t) = \psi^*(\vec{x}, -t)$

$$\Rightarrow \left(i\hbar \frac{\partial}{\partial t} + \frac{\hbar^2}{2m} \Delta \right) \psi'_T(\vec{x}, t) = 0$$

by analogy: time inversion for Dirac spinors:

$$\psi(x) \xrightarrow{T} \boxed{\psi'_T(x) = D(T) \psi^*(-\tilde{x})}$$

representation matrix for time inversion for Dirac spinors

Expectations:

$$D(T)^2 = 1, \quad D(T)^{-1} D(L_R) D(T) = D(L_R), \quad D(T)^{-2} D(M_R) D(M_A) = -D(M_R)$$

$$(i \gamma^\mu \partial_\mu - m) \psi(x) \xrightarrow{T} (i \gamma^\mu \partial_\mu - m) \psi_T(x) = 0$$

$$(i \gamma^\mu \partial_\mu - m) D(T) \psi^*(-\vec{x}) = 0 \quad | \quad D(T)^{-2}.$$

$$\left[i D(T)^{-2} \gamma^\mu D(T) \partial_\mu - m \frac{D(T)^{-2} D(T)}{=1} \right] \psi^*(-\vec{x}) = 0 \quad | \quad *$$

$$-i \left[D(T)^{-2} \gamma^\mu D(T) \right]^* \partial_\mu \psi(-\vec{x}) - m \psi(-\vec{x}) = 0$$

$$\xrightarrow{x \rightarrow -\vec{x}} \underbrace{(-i \gamma^\mu \tilde{\partial}_\mu - m)}_{= \tilde{\gamma}^\mu \partial_\mu} \psi(-\vec{x}) = 0$$

$$\left[D(T)^{-2} \gamma^\mu D(T) \right]^* = \tilde{\gamma}^\mu \Rightarrow D(T)^{-2} \gamma^\mu D(T) = (\tilde{\gamma}^\mu)^*$$

$$\gamma^\mu = \begin{pmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{pmatrix} \Rightarrow (\gamma^0)^* = \gamma^0, (\gamma^1)^* = \gamma^1, (\gamma^2)^* = -\gamma^2, (\gamma^3)^* = \gamma^3$$

$$\Rightarrow \dots \Rightarrow (\gamma^\mu)^* = (\tilde{\gamma}^\mu)^T \Rightarrow (\gamma^\mu)^\dagger = \tilde{\gamma}^\mu$$

$$\Rightarrow D(T)^{-2} \gamma^\mu D(T) = (\gamma^\mu)^T \quad \text{has to be solved for } D(T)$$

Remark: $C^{-2} \gamma^\mu C = -(\gamma^\mu)^T$

$$D(T)^{-1} \gamma^\mu D(T) = -C^{-1} \gamma^\mu C \quad | \quad C \cdot \cdot C^{-1}$$

$$C D(T)^{-2} \gamma^{\mu} D(T) C^{-1} = -\gamma^{\mu}$$

$$= (D(T) C^{-1})^{-1} [D(T) C^{-1}]^{-2} \gamma^{\mu} [D(T) C^{-1}] = -\gamma^{\mu}$$

Solution: $D(T) C^{-1} = -i \gamma^5 \Rightarrow$ you exercise

Note: Technically here all 3 discrete transformations are connected.

$$\Rightarrow D(T) = -i \gamma^5 C = \dots = i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Properties: $D(T) = D(T)^{-1} = D(T)^\dagger = -D(T)^*$

Here \dagger is not involutive in spinor space.

$$\begin{aligned} \psi_{\mp}^{\dagger}(x) &= D(T) \psi_{\mp}^{\dagger}(-\tilde{x}) = D(T) \left\{ D(T) \psi^{\dagger}(-\tilde{x}) \right\}_{x \rightarrow -\tilde{x}}^* = \underbrace{D(T) D(T)^*}_{= -D(T) D(T)^{-1}} \psi(x) \\ &= -\psi(x) \end{aligned}$$

\Rightarrow property of Dirac spinors

For you to show:

$$D(T)^{-2} D(L_{\mathbf{a}}) D(T) = D(L_{\mathbf{a}}) \neq D(L_{\mathbf{a}}) \quad \text{Expectations}$$

$$D(T)^{-2} D(M_{\mathbf{a}}) D(T) = -D(M_{\mathbf{a}}) \neq -D(M_{\mathbf{a}}) \quad \text{not fulfilled}$$

In second quantization time evolution is defined with

an anti-linear operator

$$\hat{J} (\alpha_1 \hat{O}_1 + \alpha_2 \hat{O}_2) = \alpha_1^* \hat{J}(\hat{O}_1) + \alpha_2^* \hat{J}(\hat{O}_2)$$

anti-linear

\uparrow \uparrow \uparrow \uparrow
c numbers \hat{O} and quantized operators