

Status:

field operator: $(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta) \hat{\vec{A}}(\vec{x}, t) = 0, \hat{\vec{A}}(\vec{x}, t) = \hat{\vec{A}}^\dagger(\vec{x}, t)$

$$\hat{\vec{A}}(\vec{x}, t) = \int d^3k \left\{ \hat{\vec{A}}(\vec{k}) e^{i(\vec{k} \cdot \vec{x} - \omega_{\vec{k}} t)} + \hat{\vec{A}}^\dagger(\vec{k}) e^{-i(\vec{k} \cdot \vec{x} - \omega_{\vec{k}} t)} \right\}$$

$$= N_{\vec{k}} \sum_{\lambda = \pm 1} \underbrace{\vec{E}(\vec{k}, \lambda)}_{\substack{\text{polarisation} \\ \text{vector}}} \underbrace{\hat{a}_{\vec{k}, \lambda}}_{\text{Fourier operators}}$$

helicity

+ spin $\vec{s} \parallel \vec{k}$ - spin \vec{s} antiparallel to \vec{k}

$$\vec{E}(\vec{k}, \lambda) = \begin{pmatrix} \cos \vartheta \cos \varphi - \lambda i \sin \varphi \\ \cos \vartheta \sin \varphi + \lambda i \cos \varphi \\ -\sin \vartheta \end{pmatrix}, \quad \vec{k} = k \begin{pmatrix} \sin \vartheta \cos \varphi \\ \sin \vartheta \sin \varphi \\ \cos \vartheta \end{pmatrix}$$

2.10 Properties of Polarisation Vectors:

Coulomb gauge: $\text{div } \hat{\vec{A}}(\vec{x}, t) = 0 \Leftrightarrow \vec{k} \cdot \vec{E}(\vec{k}, \lambda) = 0$

$$\vec{k} \cdot \vec{E}(\vec{k}, \lambda) = k \left\{ \sin \vartheta \cos \varphi (\cos \vartheta \cos \varphi - \lambda i \sin \varphi) \right. \\ \left. + \sin \vartheta \sin \varphi (\cos \vartheta \sin \varphi + \lambda i \cos \varphi) - \sin \vartheta \cos \vartheta \right\} \stackrel{\text{transversality condition}}{=} 0 \checkmark$$

orthonormality relation:

$$\vec{E}^*(\vec{k}, \lambda) \cdot \vec{E}(\vec{k}, \lambda') \stackrel{\text{exercise}}{=} \delta_{\lambda \lambda'}; \quad \lambda, \lambda' = \pm 1$$

2) Behaviour under inversion: $\vec{k} \Rightarrow -\vec{k}$

$$\left. \begin{array}{l} \varphi \rightarrow \varphi + \pi: \quad \sin \varphi \rightarrow -\sin \varphi, \quad \cos \varphi \rightarrow -\cos \varphi \\ \vartheta \rightarrow \vartheta - \pi: \quad \sin \vartheta \rightarrow \sin \vartheta, \quad \cos \vartheta \rightarrow -\cos \vartheta \\ \vec{k} \rightarrow -\vec{k} \quad (\text{spherical coordinates}) \end{array} \right\} \begin{array}{l} \vec{E}(-\vec{k}, \lambda) \stackrel{\text{exercise}}{=} \vec{E}(\vec{k}, -\lambda) \\ = E^*(\vec{k}, \lambda) \end{array}$$

$$\hat{\vec{A}}(\vec{x}, t) = \sum_{\lambda=\pm 1} \int d^3k N_{\vec{k}} \left\{ \vec{E}(\vec{k}, \lambda) e^{i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda} + \vec{E}^*(\vec{k}, \lambda) e^{-i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda}^{\dagger} \right\}$$

What is the physical meaning of these Fourier operators?

2.11 Fourier Operators:

$$\hat{\vec{\Pi}}(\vec{x}, t) = \epsilon_0 \frac{\partial \hat{\vec{A}}(\vec{x}, t)}{\partial t} = \epsilon_0 \sum_{\lambda=\pm 1} \int d^3k N_{\vec{k}} \left\{ -i\omega_{\vec{k}} \vec{E}(\vec{k}, t) e^{i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda} + i\omega_{\vec{k}} \vec{E}^*(\vec{k}, t) e^{-i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda}^{\dagger} \right\}$$

Note: $[\hat{A}_i(\vec{x}, t), \hat{\Pi}_j(\vec{x}', t)] = i\hbar \delta_{ij}^T(\vec{x} - \vec{x}')$; $[\hat{A}_i(\vec{x}, t), \hat{A}_j(\vec{x}', t)] = 0 = [\hat{\Pi}_i(\vec{x}, t), \hat{\Pi}_j(\vec{x}', t)]$

transversal delta function

exercise

$$\hat{a}_{\vec{k}, \lambda} = \frac{1}{2(2\pi)^3 N_{\vec{k}}} \int d^3x \vec{E}^*(\vec{k}, \lambda) \cdot e^{-i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \left\{ \hat{\vec{A}}(\vec{x}, t) + \frac{i}{\epsilon_0 \omega_{\vec{k}}} \hat{\vec{\Pi}}(\vec{x}, t) \right\}$$

$$\hat{a}_{\vec{k}, \lambda}^{\dagger} = \frac{1}{2(2\pi)^3 N_{\vec{k}}} \int d^3x \vec{E}(\vec{k}, \lambda) \cdot e^{+i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \left\{ \hat{\vec{A}}(\vec{x}, t) - \frac{i}{\epsilon_0 \omega_{\vec{k}}} \hat{\vec{\Pi}}(\vec{x}, t) \right\}$$

$$[\hat{a}_{\vec{k}, \lambda}, \hat{a}_{\vec{k}', \lambda'}] = \dots = 0 = [\hat{a}_{\vec{k}, \lambda}^{\dagger}, \hat{a}_{\vec{k}', \lambda'}^{\dagger}]$$

$$[\hat{a}_{\vec{k}, \lambda}, \hat{a}_{\vec{k}', \lambda'}]_- = -\dots = \frac{\hbar}{2(2\pi)^3 \epsilon_0 \omega_{\vec{k}} N_{\vec{k}}^2} \delta(\vec{k} - \vec{k}') \delta_{\lambda\lambda'}$$

↑
transversal delta function

represent annihilation and creation operators of bosonic

particles, which are characterized by \vec{k}, λ

\Rightarrow photons

$$\stackrel{!}{=} 1 \Rightarrow N_{\vec{k}} = \sqrt{\frac{\hbar}{2(2\pi)^3 \epsilon_0 \omega_{\vec{k}}}}$$

2.12 Energy:

$$\hat{\vec{A}}(\vec{x}, t) = \sum_{\lambda=\pm 1} \int d^3k \sqrt{\frac{\hbar}{2(2\pi)^3 \epsilon_0 \omega_{\vec{k}}}} \left\{ \vec{\epsilon}(\vec{k}, \lambda) e^{i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda} + \vec{\epsilon}^*(\vec{k}, \lambda) e^{-i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda}^{\dagger} \right\}$$

$$\hat{\vec{\pi}}(\vec{x}, t) = \sum_{\lambda=\pm 1} \int d^3k \sqrt{\frac{\hbar \omega_{\vec{k}} \epsilon_0}{2(2\pi)^3}} \left\{ -i \vec{\epsilon}(\vec{k}, \lambda) e^{i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda} + i \vec{\epsilon}^*(\vec{k}, \lambda) e^{-i(\vec{k}\vec{x} - \omega_{\vec{k}}t)} \hat{a}_{\vec{k}, \lambda}^{\dagger} \right\}$$

$$\hat{H} = \frac{1}{2} \int d^3x \left\{ \underbrace{\frac{1}{\epsilon_0} \hat{\vec{\pi}}(\vec{x}, t) \cdot \hat{\vec{\pi}}(\vec{x}, t)}_{\hat{=} \text{ electric energy}} + \underbrace{\frac{1}{\mu_0} \partial_{\lambda} \hat{\vec{A}}(\vec{x}, t) \cdot \partial_{\lambda} \hat{\vec{A}}(\vec{x}, t)}_{\hat{=} \text{ magnetic energy}} \right\}$$

exercise

$$= \dots = \frac{1}{2} \sum_{\lambda=\pm 1} \int d^3k \underbrace{\hbar \omega_{\vec{k}}}_{=c|\vec{k}|} \left\{ \hat{a}_{\vec{k}, \lambda}^{\dagger} \hat{a}_{\vec{k}, \lambda} + \hat{a}_{\vec{k}, \lambda} \hat{a}_{\vec{k}, \lambda}^{\dagger} \right\}$$

- Time independent despite of having time dependent $\hat{\vec{A}}(\vec{x}, t), \hat{\vec{\pi}}(\vec{x}, t)$ due to conservation of energy

- Electromagnetic field consists of independent harmonic oscillators: each energy quantum $\hbar \omega_{\vec{k}}$ is doubly degenerate due to $\lambda = \pm 1$

vacuum state: $\hat{a}_{\vec{k}, \lambda} |0\rangle = 0 = \langle 0 | \hat{a}_{\vec{k}, \lambda}^\dagger$ divergent

vacuum energy: $\langle 0 | \hat{H} | 0 \rangle = \frac{1}{2} \sum_{\lambda=\pm 1} \int d^3x \underbrace{\epsilon_0 \hbar \omega_{\vec{k}}}_{\infty} \underbrace{\delta(\vec{0})}_{\infty} = \infty$

renormalized Hamiltonian operator:

$$\begin{aligned} \text{:} \hat{H} \text{:} &= \hat{H} - \langle 0 | \hat{H} | 0 \rangle = \frac{1}{2} \sum_{\lambda=\pm 1} \int d^3x \epsilon_0 \hbar \omega_{\vec{k}} (\hat{a}_{\vec{k}, \lambda}^\dagger \hat{a}_{\vec{k}, \lambda} + \hat{a}_{\vec{k}, \lambda} \hat{a}_{\vec{k}, \lambda}^\dagger) - \int d^3x \epsilon_0 \hbar \omega_{\vec{k}} \delta(\vec{0}) \\ &= \sum_{\lambda=\pm 1} \int d^3x \epsilon_0 \hbar \omega_{\vec{k}} \underbrace{\hat{a}_{\vec{k}, \lambda}^\dagger \hat{a}_{\vec{k}, \lambda}}_{= \hat{n}_{\vec{k}, \lambda} \text{ photon number operator}} = \hat{a}_{\vec{k}, \lambda}^\dagger \hat{a}_{\vec{k}, \lambda} + \int d^3x \epsilon_0 \hbar \omega_{\vec{k}} \delta(\vec{0}) \end{aligned}$$

normal ordering

→ creation operators to the left, annihilation operators to the right

2.13 Momentum:

Emmy Noether's theorem: $\vec{P} = \int d^3x \frac{\vec{S}(\vec{x}, t)}{c^2}$

Poynting vector: $\vec{S}(\vec{x}, t) = \frac{1}{\mu_0} \underbrace{\vec{E}(\vec{x}, t)}_{= -\vec{\nabla} \phi} \times \underbrace{\vec{B}(\vec{x}, t)}_{\text{rot } \vec{A}}$

$$\Rightarrow \hat{\vec{P}} = \int d^3x [\vec{\nabla} \times \hat{\vec{A}}(\vec{x}, t)] \times \hat{\vec{\nabla}} \phi(\vec{x}, t)$$

→ plane wave decomposition

$$\hat{\vec{P}} = \dots = \sum_{\lambda=\pm 1} \int d^3x \frac{\hbar \vec{k}}{2} \{ \hat{a}_{\vec{k}, \lambda}^\dagger \hat{a}_{\vec{k}, \lambda} + \hat{a}_{\vec{k}, \lambda} \hat{a}_{\vec{k}, \lambda}^\dagger \}$$

$$\langle 0 | \hat{\vec{P}} | 0 \rangle = \underbrace{\int d^3x \hbar \vec{k}}_{=0} \delta(\vec{0}) = 0$$

$$:\hat{\vec{p}}: = \hat{\vec{p}} = \sum_{\lambda=\pm 1} \int d^3k \frac{1}{k} \vec{k} \hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda}$$

2.14 Spin Angular Momentum:

Noether theorem:

$$\hat{\vec{S}} = \int d^3x \hat{\vec{A}}(\vec{x}, t) \times \hat{\vec{\pi}}(\vec{x}, t)$$

$$= \vec{L} = \vec{x} \times \vec{p}$$

$$= \dots = \sum_{\lambda=\pm 1} \int d^3k \lambda \frac{1}{2} \frac{\vec{k}}{|\vec{k}|} (\hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda} + \hat{a}_{\vec{k},\lambda} \hat{a}_{\vec{k},\lambda}^\dagger)$$

$$\langle 0 | \hat{\vec{S}} | 0 \rangle = \underbrace{\sum_{\lambda=\pm 1} \lambda}_{=0} \left(\int d^3k \frac{\vec{k}}{|\vec{k}|} \right) \underbrace{\delta(0)}_{=0} = 0$$

$$:\hat{\vec{S}}: = \hat{\vec{S}} = \sum_{\lambda=\pm 1} \int d^3k \lambda \frac{1}{2} \frac{\vec{k}}{|\vec{k}|} \hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda}$$

helicity

photon has spin 1

2.15 Fock Basis:

first quantized harmonic oscillator
(Appendix F)

second quantized electromagnetic field

lowering / raising operators \hat{a} , \hat{a}^\dagger

annihilation / creation operators

$\hat{a}_{\vec{k},\lambda}$, $\hat{a}_{\vec{k},\lambda}^\dagger$

$$[\hat{a}, \hat{a}]_- = [\hat{a}^\dagger, \hat{a}^\dagger]_- = 0$$

$$[\hat{a}_{\vec{k},\lambda}, \hat{a}_{\vec{k}',\lambda'}]_- = 0 = [\hat{a}_{\vec{k},\lambda}^\dagger, \hat{a}_{\vec{k}',\lambda'}^\dagger]_-$$

$$[\hat{a}, \hat{a}^\dagger]_- = 1$$

ground state $|0\rangle$

$$\hat{a}|0\rangle = 0 = \langle 0|\hat{a}^\dagger$$

eigenstates of \hat{H}

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle$$

orthonormal:

$$\langle n | n' \rangle = \delta_{n, n'}$$

completeness:

$$\sum_{n=0}^{\infty} |n\rangle \langle n| = 1$$

first quantized Hilbert space

$$\mathcal{H} = \text{Span} \{ |0\rangle, |1\rangle, |2\rangle, \dots \}$$

$$[\hat{a}_{\vec{\mu}, \lambda}, \hat{a}_{\vec{\mu}', \lambda'}^\dagger]_- = \delta(\vec{\mu} - \vec{\mu}') \delta_{\lambda \lambda'}$$

vacuum state $|0\rangle$

$$\hat{a}_{\vec{\mu}, \lambda} |0\rangle = 0 = \langle 0| \hat{a}_{\vec{\mu}, \lambda}^\dagger$$

eigenstates of \hat{H}

$$|\{n_{\vec{\mu}, \lambda}\}\rangle = \prod_{\vec{\mu}, \lambda} \frac{1}{\sqrt{n_{\vec{\mu}, \lambda}!}} (\hat{a}_{\vec{\mu}, \lambda}^\dagger)^{n_{\vec{\mu}, \lambda}} |0\rangle$$

\uparrow set of all
 \uparrow photon numbers
 Fock states

orthonormality

$$\langle \{n_{\vec{\mu}, \lambda}\} | \{n'_{\vec{\mu}, \lambda}\} \rangle = \prod_{\vec{\mu}, \lambda} \delta_{n_{\vec{\mu}, \lambda}, n'_{\vec{\mu}, \lambda}}$$

completeness

$$\sum_{\lambda=\pm 1} \int d^3q \sum_{n_{\vec{\mu}, \lambda}=0}^{\infty} |\{n_{\vec{\mu}, \lambda}\}\rangle \langle \{n_{\vec{\mu}, \lambda}\}| = 1$$

second quantized Fock space

$$\mathcal{F} = \text{Span} \{ |0\rangle, |1_{\vec{\mu}, \lambda}\rangle, |1_{\vec{\mu}, \lambda}; 1_{\vec{\mu}', \lambda'}\rangle, |2_{\vec{\mu}, \lambda}\rangle, \dots \}$$

first quantized Hamilton operator second quantized Hamilton operator

$$\hat{H} = \hbar\omega \left(\hat{n} + \frac{1}{2} \right)$$

$\hat{n} = \hat{a}^\dagger \hat{a}$

$$\hat{H} = \sum_{\lambda=\pm 1} \int d^3k \, \hbar\omega_{\vec{k}} \underbrace{\hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda}}_{= \hat{n}_{\vec{k},\lambda}}$$