

Quantum Optics

Problem Sheet 11

Problem 24: Atom Interferometer

(10 points)

In this exercise we derive a formalism for atomic interferometry, which is based on the optical Bloch equations with vanishing detuning Δ .

a) Solve the optical Bloch equations for the Bloch vector $\mathbf{s}(t)$ in the special case $\Delta = 0$, i.e.

$$\dot{\mathbf{s}}(t) = -\Omega_R \mathbf{e}_1 \times \mathbf{s}(t), \quad (1)$$

where Ω_R denotes the Rabi frequency. Write down an explicit expression for the time-evolution operator $\hat{U}(t)$.

b) Consider now as the initial state that the atom is in the ground state. What is the final state for times $\Omega_R t = \pi/2$ and $\Omega_R t = \pi$? How does the operator for measurement of the population imbalance s_3 read, and how can you calculate the actual population of ground and excited state from this operator?

c) Can you construct a beam splitter for atoms with your results from a) and b)?

d) Describe how you can build up an atomic Mach-Zehnder interferometer and calculate the evolution of the atom beam through the interferometer.

e) Explain qualitatively how the results change in presence of a detuning Δ .

Problem 25: Bloch Equations and Decay

(14 points)

Real experiments with atomic beams suffer from several decay mechanisms like population decay via spontaneous emission, which is modelled via the T_1 -time, and decoherence of the atomic phases through varying detuning Δ , which is described via the T_2 -time. In this exercise, we will simulate these processes stochastically.

a) Implement the optical Bloch equations with detuning Δ

$$\dot{\mathbf{s}}(t) = - \begin{pmatrix} \Omega_R \\ 0 \\ \Delta \end{pmatrix} \times \mathbf{s}(t) \quad (2)$$

numerically.

In the following we add peu à peu the stochastic processes.

b) At first we focus on the spontaneous decay, which means that at some random time the excited states decay into the ground state.

The spontaneous emission is included numerically by drawing a uniformly distributed real number r from the interval $[0, 1]$. If $r < \Delta t \gamma$, where Δt is the time step and γ the decay rate, we set $s_3 = -1$, otherwise we perform a normal time step. Show numerically that in the special case $\Omega_R = 0$ this procedure yields in the average of many realisations

an exponential damping $e^{-\gamma t}$ for the excited state population. Which initial condition makes the most sense? What happens for finite frequency Ω_R ?

c) We come now to the decoherence, which means that the detuning Δ is not constant but varies from atom to atom. We assume here a Gaussian distributed detuning, which can arise, for instance, in atomic beams through the Doppler effect.

Use your implementation from **a)** with the initial state $\mathbf{s}_0 = (0, 1, 0)^T$ and a Gaussian distributed detuning Δ with $\langle \Delta \rangle = 0$ and $\langle \Delta^2 \rangle \sim T$, where T is the temperature of the surrounding introducing the Doppler shift. Show that this leads in the ensemble average to a decay of the Bloch vector.

d) We now combine the effects from **b)** and **c)** and implement a typical experimental sequence, namely the spin-echo sequence. At the beginning, the atom is considered to be in the ground state. The sequence consists of a $\pi/2$ -pulse followed by a waiting time τ , then a π -pulse again followed by a waiting time τ and finally a $\pi/2$ -pulse again. What do you observe for the Bloch vector for different waiting times τ and how do you explain your findings?

e) Implement your Mach-Zehnder interferometer from **d)** and add spontaneous decay and decoherence. How do the results change?

Drop the solutions in the post box on the 5th floor of building 46 or, in case of illness/quarantine, send them via email to jkrauss@rhrk.uni-kl.de until July 7 at 12.00.